

ma01ag

Constructs matrices that represent the Beltrami linear system.

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1.1 gauge conditions

- In the v -th annulus, bounded by the $(v - 1)$ -th and v -th interfaces, a general covariant representation of the magnetic vector-potential is written

$$\bar{\mathbf{A}} = \bar{A}_s \nabla s + \bar{A}_\theta \nabla \theta + \bar{A}_\zeta \nabla \zeta. \quad (1)$$

- To this add $\nabla g(s, \theta, \zeta)$, where g satisfies

$$\begin{aligned} \partial_s g(s, \theta, \zeta) &= -\bar{A}_s(s, \theta, \zeta) \\ \partial_\theta g(-1, \theta, \zeta) &= -\bar{A}_\theta(-1, \theta, \zeta) \\ \partial_\zeta g(-1, 0, \zeta) &= -\bar{A}_\zeta(-1, 0, \zeta). \end{aligned} \quad (2)$$

- Then $\mathbf{A} = \bar{\mathbf{A}} + \nabla g$ is given by $\mathbf{A} = A_\theta \nabla \theta + A_\zeta \nabla \zeta$ with

$$A_\theta(-1, \theta, \zeta) = 0 \quad (3)$$

$$A_\zeta(-1, 0, \zeta) = 0 \quad (4)$$

- This specifies the gauge: to see this, notice that no gauge term can be added without violating the conditions in Eq.(3) or Eq.(4).
- Note that the gauge employed in each volume is distinct.

1.2 boundary conditions

- The magnetic field is $\sqrt{g} \mathbf{B} = (\partial_\theta A_\zeta - \partial_\zeta A_\theta) \mathbf{e}_s - \partial_s A_\zeta \mathbf{e}_\theta + \partial_s A_\theta \mathbf{e}_\zeta$.
- In the annular volumes, the condition that the field is tangential to the inner interface gives $\partial_\theta A_\zeta - \partial_\zeta A_\theta = 0$. With the above gauge condition on A_θ given in Eq.(3), this gives $\partial_\theta A_\zeta = 0$, which with Eq.(4) gives

$$A_\zeta(-1, \theta, \zeta) = 0. \quad (5)$$

- The condition at the outer interface is that the field is $\sqrt{g} \mathbf{B} \cdot \nabla s = \partial_\theta A_\zeta - \partial_\zeta A_\theta = b$, where b is supplied by the user. For each of the plasma regions, $b = 0$. For the vacuum region, generally $b \neq 0$.

1.3 enclosed fluxes

- In the plasma regions, the enclosed fluxes must be constrained.
- The toroidal and poloidal fluxes enclosed in each volume are determined using

$$\int_S \mathbf{B} \cdot d\mathbf{s} = \int_{\partial S} \mathbf{A} \cdot d\mathbf{l}. \quad (6)$$

1.4 Fourier-Chebyshev representation

- The components of the vector potential are

$$A_\theta(s, \theta, \zeta) = \sum_{i,l} \textcolor{red}{A}_{\theta,e,i,l} T_l(s) \cos \alpha_i + \sum_{i,l} \textcolor{orange}{A}_{\theta,o,i,l} T_l(s) \sin \alpha_i, \quad (7)$$

$$A_\zeta(s, \theta, \zeta) = \sum_{i,l} \textcolor{blue}{A}_{\zeta,e,i,l} T_l(s) \cos \alpha_i + \sum_{i,l} \textcolor{blue}{A}_{\zeta,o,i,l} T_l(s) \sin \alpha_i, \quad (8)$$

where $\alpha_j \equiv m_j \theta - n_j \zeta$.

1.5 constrained energy functional

- The constrained energy functional in each volume is

$$\begin{aligned} \mathcal{F}_l &\equiv \int \mathbf{B} \cdot \mathbf{B} dv - \frac{\mu}{2} \left[\int \mathbf{A} \cdot \mathbf{B} dv - K_l \right] \\ &+ \sum_i a_i \begin{bmatrix} \sum_l \textcolor{red}{A}_{\theta,e,i,l} T_l(-1) - 0 \\ \sum_l \textcolor{blue}{A}_{\zeta,e,i,l} T_l(-1) - 0 \\ \sum_l \textcolor{orange}{A}_{\theta,o,i,l} T_l(-1) - 0 \\ \sum_l \textcolor{blue}{A}_{\zeta,o,i,l} T_l(-1) - 0 \end{bmatrix} \\ &+ \sum_i b_i \begin{bmatrix} \sum_l (-m_i \textcolor{blue}{A}_{\zeta,e,i,l} - n_i \textcolor{red}{A}_{\theta,e,i,l}) T_l(+1) - b_{s,i} \\ \sum_l (+m_i \textcolor{blue}{A}_{\zeta,o,i,l} + n_i \textcolor{orange}{A}_{\theta,o,i,l}) T_l(+1) - b_{c,i} \end{bmatrix} \\ &+ \sum_i e_i \begin{bmatrix} \sum_l \textcolor{red}{A}_{\theta,e,1,l} T_l(+1) - \Delta \psi_t \\ \sum_l \textcolor{blue}{A}_{\zeta,e,1,l} T_l(+1) - \Delta \psi_p \end{bmatrix} \\ &+ \alpha \begin{bmatrix} \sum_l \textcolor{red}{A}_{\theta,e,1,l} T_l(+1) - \Delta \psi_t \\ \sum_l \textcolor{blue}{A}_{\zeta,e,1,l} T_l(+1) - \Delta \psi_p \end{bmatrix} \end{aligned} \quad (9)$$

where

- a_i, b_i, c_i and d_i are Lagrange multipliers used to enforce the combined gauge and interface boundary condition on the inner interface,
 - e_i and f_i are Lagrange multipliers used to enforce the interface boundary condition on the outer interface, namely $\sqrt{g} \mathbf{B} \cdot \nabla s = b$; and
 - α and β are Lagrange multipliers used to enforce the constraints on the enclosed fluxes.
- In each plasma volume the boundary condition on the outer interface is $b = 0$.
 - In the vacuum volume (only for free-boundary), we may set $\mu = 0$.

1.6 energy and helicity integrands

1. The integrands are

$$\begin{aligned}
\sqrt{g} \mathbf{B} \cdot \mathbf{B} = & + (m_i A_{\zeta,o,i,l} + n_i A_{\theta,o,i,l}) (m_j A_{\zeta,o,j,p} + n_j A_{\theta,o,j,p}) T_l T_p \bar{g}_{ss} \cos \alpha_i \cos \alpha_j \\
& - 2 (m_i A_{\zeta,o,i,l} + n_i A_{\theta,o,i,l}) (m_j A_{\zeta,e,j,p} + n_j A_{\theta,e,j,p}) T_l T_p \bar{g}_{ss} \cos \alpha_i \sin \alpha_j \\
& + (m_i A_{\zeta,e,i,l} + n_i A_{\theta,e,i,l}) (m_j A_{\zeta,e,j,p} + n_j A_{\theta,e,j,p}) T_l T_p \bar{g}_{ss} \sin \alpha_i \sin \alpha_j \\
& - 2 (m_i A_{\zeta,o,i,l} + n_i A_{\theta,o,i,l}) A_{\zeta,e,j,p} T_l T'_p \bar{g}_{s\theta} \cos \alpha_i \cos \alpha_j \\
& - 2 (m_i A_{\zeta,o,i,l} + n_i A_{\theta,o,i,l}) A_{\zeta,o,j,p} T_l T'_p \bar{g}_{s\theta} \cos \alpha_i \sin \alpha_j \\
& + 2 (m_i A_{\zeta,e,i,l} + n_i A_{\theta,e,i,l}) A_{\zeta,e,j,p} T_l T'_p \bar{g}_{s\theta} \sin \alpha_i \cos \alpha_j \\
& + 2 (m_i A_{\zeta,e,i,l} + n_i A_{\theta,e,i,l}) A_{\zeta,o,j,p} T_l T'_p \bar{g}_{s\theta} \sin \alpha_i \sin \alpha_j \\
& + 2 (m_i A_{\zeta,o,i,l} + n_i A_{\theta,o,i,l}) A_{\theta,e,j,p} T_l T'_p \bar{g}_{s\zeta} \cos \alpha_i \cos \alpha_j \\
& + 2 (m_i A_{\zeta,o,i,l} + n_i A_{\theta,o,i,l}) A_{\theta,o,j,p} T_l T'_p \bar{g}_{s\zeta} \cos \alpha_i \sin \alpha_j \\
& - 2 (m_i A_{\zeta,e,i,l} + n_i A_{\theta,e,i,l}) A_{\theta,e,j,p} T_l T'_p \bar{g}_{s\zeta} \sin \alpha_i \cos \alpha_j \\
& - 2 (m_i A_{\zeta,e,i,l} + n_i A_{\theta,e,i,l}) A_{\theta,o,j,p} T_l T'_p \bar{g}_{s\zeta} \sin \alpha_i \sin \alpha_j \\
& + A_{\zeta,e,i,l} A_{\zeta,e,j,p} T'_l T'_p \bar{g}_{\theta\theta} \cos \alpha_i \cos \alpha_j \\
& + 2 A_{\zeta,e,i,l} A_{\zeta,o,j,p} T'_l T'_p \bar{g}_{\theta\theta} \cos \alpha_i \sin \alpha_j \\
& + A_{\zeta,o,i,l} A_{\zeta,o,j,p} T'_l T'_p \bar{g}_{\theta\theta} \sin \alpha_i \sin \alpha_j \\
& - 2 A_{\zeta,e,i,l} A_{\theta,e,j,p} T'_l T'_p \bar{g}_{\theta\zeta} \cos \alpha_i \cos \alpha_j \\
& - 2 A_{\zeta,e,i,l} A_{\theta,o,j,p} T'_l T'_p \bar{g}_{\theta\zeta} \cos \alpha_i \sin \alpha_j \\
& - 2 A_{\zeta,o,i,l} A_{\theta,e,j,p} T'_l T'_p \bar{g}_{\theta\zeta} \sin \alpha_i \cos \alpha_j \\
& - 2 A_{\zeta,o,i,l} A_{\theta,o,j,p} T'_l T'_p \bar{g}_{\theta\zeta} \sin \alpha_i \sin \alpha_j \\
& + A_{\theta,e,i,l} A_{\theta,e,j,p} T'_l T'_p \bar{g}_{\zeta\zeta} \cos \alpha_i \cos \alpha_j \\
& + 2 A_{\theta,e,i,l} A_{\theta,o,j,p} T'_l T'_p \bar{g}_{\zeta\zeta} \cos \alpha_i \sin \alpha_j \\
& + A_{\theta,o,i,l} A_{\zeta,e,j,p} T'_l T'_p \bar{g}_{\zeta\zeta} \sin \alpha_i \cos \alpha_j \\
& + A_{\theta,o,i,l} A_{\zeta,o,j,p} T'_l T'_p \bar{g}_{\zeta\zeta} \sin \alpha_i \sin \alpha_j
\end{aligned} \tag{10}$$

$$\begin{aligned}
\sqrt{g} \mathbf{A} \cdot \mathbf{B} = & - A_{\zeta,e,i,l} A_{\theta,e,j,p} T'_l T_p \cos \alpha_i \cos \alpha_j \\
& - A_{\zeta,e,i,l} A_{\theta,o,j,p} T'_l T_p \cos \alpha_i \sin \alpha_j \\
& - A_{\zeta,o,i,l} A_{\theta,e,j,p} T'_l T_p \sin \alpha_i \cos \alpha_j \\
& - A_{\zeta,o,i,l} A_{\theta,o,j,p} T'_l T_p \sin \alpha_i \sin \alpha_j \\
& + A_{\theta,e,i,l} A_{\zeta,e,j,p} T'_l T_p \cos \alpha_i \cos \alpha_j \\
& + A_{\theta,e,i,l} A_{\zeta,o,j,p} T'_l T_p \cos \alpha_i \sin \alpha_j \\
& + A_{\theta,o,i,l} A_{\zeta,e,j,p} T'_l T_p \sin \alpha_i \cos \alpha_j \\
& + A_{\theta,o,i,l} A_{\zeta,o,j,p} T'_l T_p \sin \alpha_i \sin \alpha_j
\end{aligned} \tag{11}$$

1.7 first derivatives of energy and helicity integrands with respect to $A_{\theta,e,i,l}$ and $A_{\theta,o,i,l}$

1. The first derivatives with respect to $A_{\theta,e,i,l}$ and $A_{\theta,o,i,l}$ are

$$\begin{aligned} \frac{\partial}{\partial A_{\theta,e,i,l}} \int dv \mathbf{B} \cdot \mathbf{B} = & \\ -2n_i(m_j A_{\zeta,o,j,p} + n_j A_{\theta,o,j,p}) & \int ds T_p T_l \oint d\theta d\zeta \bar{g}_{ss} \cos \alpha_j \sin \alpha_i \\ +2n_i(m_j A_{\zeta,e,j,p} + n_j A_{\theta,e,j,p}) & \int ds T_l T_p \oint d\theta d\zeta \bar{g}_{ss} \sin \alpha_i \sin \alpha_j \\ +2n_i A_{\zeta,e,j,p} & \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\theta} \sin \alpha_i \cos \alpha_j \\ +2n_i A_{\zeta,o,j,p} & \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\theta} \sin \alpha_i \sin \alpha_j \\ +2(m_j A_{\zeta,o,j,p} + n_j A_{\theta,o,j,p}) & \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\zeta} \cos \alpha_j \cos \alpha_i \\ -2n_i A_{\theta,e,j,p} & \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\zeta} \sin \alpha_i \cos \alpha_j \\ -2(m_j A_{\zeta,e,j,p} + n_j A_{\theta,e,j,p}) & \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\zeta} \sin \alpha_j \cos \alpha_i \\ -2A_{\theta,o,j,p} & \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\zeta} \sin \alpha_i \sin \alpha_j \\ -2A_{\zeta,e,j,p} & \int ds T'_p T'_l \oint d\theta d\zeta \bar{g}_{\theta\zeta} \cos \alpha_j \cos \alpha_i \\ -2A_{\zeta,o,j,p} & \int ds T'_p T'_l \oint d\theta d\zeta \bar{g}_{\theta\zeta} \sin \alpha_j \cos \alpha_i \\ +2A_{\theta,e,j,p} & \int ds T'_l T'_p \oint d\theta d\zeta \bar{g}_{\zeta\zeta} \cos \alpha_i \cos \alpha_j \\ +2A_{\theta,o,j,p} & \int ds T'_l T'_p \oint d\theta d\zeta \bar{g}_{\zeta\zeta} \cos \alpha_i \sin \alpha_j \end{aligned} \quad (12)$$

$$\begin{aligned} \frac{\partial}{\partial A_{\theta,e,i,l}} \int dv \mathbf{A} \cdot \mathbf{B} = & \\ -A_{\zeta,e,j,p} \int ds T'_p T_l \oint d\theta d\zeta & \cos \alpha_j \cos \alpha_i \\ -A_{\zeta,o,j,p} \int ds T'_p T_l \oint d\theta d\zeta & \sin \alpha_j \cos \alpha_i \\ +A_{\zeta,e,j,p} \int ds T'_l T_p \oint d\theta d\zeta & \cos \alpha_i \cos \alpha_j \\ +A_{\zeta,o,j,p} \int ds T'_l T_p \oint d\theta d\zeta & \cos \alpha_i \sin \alpha_j \end{aligned} \quad (13)$$

$$\begin{aligned} \frac{\partial}{\partial A_{\theta,o,i,l}} \int dv \mathbf{B} \cdot \mathbf{B} = & \\ +2n_i(m_j A_{\zeta,o,j,p} + n_j A_{\theta,o,j,p}) & \int ds T_l T_p \oint d\theta d\zeta \bar{g}_{ss} \cos \alpha_i \cos \alpha_j \\ -2n_i(m_j A_{\zeta,e,j,p} + n_j A_{\theta,e,j,p}) & \int ds T_l T_p \oint d\theta d\zeta \bar{g}_{ss} \cos \alpha_i \sin \alpha_j \\ -2n_i A_{\zeta,e,j,p} & \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\theta} \cos \alpha_i \cos \alpha_j \\ -2n_i A_{\zeta,o,j,p} & \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\theta} \cos \alpha_i \sin \alpha_j \\ +2n_i A_{\theta,e,j,p} & \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\zeta} \cos \alpha_i \cos \alpha_j \\ +2n_i A_{\theta,o,j,p} & \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\zeta} \cos \alpha_i \sin \alpha_j \\ +2(m_j A_{\zeta,o,j,p} + n_j A_{\theta,o,j,p}) & \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\zeta} \cos \alpha_j \sin \alpha_i \\ -2(m_j A_{\zeta,e,j,p} + n_j A_{\theta,e,j,p}) & \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\zeta} \sin \alpha_j \sin \alpha_i \\ -2A_{\zeta,e,j,p} & \int ds T'_p T'_l \oint d\theta d\zeta \bar{g}_{\theta\zeta} \cos \alpha_j \sin \alpha_i \\ -2A_{\zeta,o,j,p} & \int ds T'_p T'_l \oint d\theta d\zeta \bar{g}_{\theta\zeta} \sin \alpha_j \sin \alpha_i \\ +2A_{\theta,e,j,p} & \int ds T'_p T'_l \oint d\theta d\zeta \bar{g}_{\zeta\zeta} \cos \alpha_j \sin \alpha_i \\ +2A_{\theta,o,j,p} & \int ds T'_l T'_p \oint d\theta d\zeta \bar{g}_{\zeta\zeta} \sin \alpha_i \sin \alpha_j \end{aligned} \quad (14)$$

$$\begin{aligned} \frac{\partial}{\partial A_{\theta,o,i,l}} \int dv \mathbf{A} \cdot \mathbf{B} = & \\ -A_{\zeta,e,j,p} \int ds T'_p T_l \oint d\theta d\zeta & \cos \alpha_j \sin \alpha_i \\ -A_{\zeta,o,j,p} \int ds T'_p T_l \oint d\theta d\zeta & \sin \alpha_j \sin \alpha_i \\ +A_{\zeta,e,j,p} \int ds T'_l T_p \oint d\theta d\zeta & \sin \alpha_i \cos \alpha_j \\ +A_{\zeta,o,j,p} \int ds T'_l T_p \oint d\theta d\zeta & \sin \alpha_i \sin \alpha_j \end{aligned} \quad (15)$$

1.8 first derivatives of energy and helicity integrands with respect to $A_{\zeta,e,i,l}$ and $A_{\zeta,o,i,l}$

1. The first derivatives with respect to $A_{\zeta,e,i,l}$ and $A_{\zeta,o,i,l}$ are

$$\begin{aligned} \frac{\partial}{\partial A_{\zeta,e,i,l}} \int dv \mathbf{B} \cdot \mathbf{B} = & \\ -2m_i(m_j A_{\zeta,o,j,p} + n_j A_{\theta,o,j,p}) & \int ds T_p T_l \oint d\theta d\zeta \bar{g}_{ss} \cos \alpha_j \sin \alpha_i \\ +2m_i(m_j A_{\zeta,e,j,p} + n_j A_{\theta,e,j,p}) & \int ds T_l T_p \oint d\theta d\zeta \bar{g}_{ss} \sin \alpha_i \sin \alpha_j \\ -2(m_j A_{\zeta,o,j,p} + n_j A_{\theta,o,j,p}) & \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\theta} \cos \alpha_j \cos \alpha_i \\ +2m_i A_{\zeta,e,j,p} & \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\theta} \sin \alpha_i \cos \alpha_j \\ +2(m_j A_{\zeta,e,j,p} + n_j A_{\theta,e,j,p}) & \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\theta} \sin \alpha_j \cos \alpha_i \\ +2m_i A_{\zeta,o,j,p} & \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\theta} \sin \alpha_i \sin \alpha_j \\ -2m_i A_{\theta,e,j,p} & \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\zeta} \sin \alpha_i \cos \alpha_j \\ -2m_i A_{\theta,o,j,p} & \int ds T'_l T_p \oint d\theta d\zeta \bar{g}_{\theta\theta} \cos \alpha_i \cos \alpha_j \\ +2A_{\zeta,e,j,p} & \int ds T'_l T'_p \oint d\theta d\zeta \bar{g}_{\theta\theta} \cos \alpha_i \sin \alpha_j \\ +2A_{\zeta,o,j,p} & \int ds T'_l T'_p \oint d\theta d\zeta \bar{g}_{\theta\zeta} \cos \alpha_i \cos \alpha_j \\ -2A_{\theta,e,j,p} & \int ds T'_l T'_p \oint d\theta d\zeta \bar{g}_{\theta\zeta} \cos \alpha_i \sin \alpha_j \\ -2A_{\theta,o,j,p} & \int ds T'_l T'_p \oint d\theta d\zeta \bar{g}_{\theta\zeta} \sin \alpha_i \sin \alpha_j \end{aligned} \quad (16)$$

$$\begin{aligned} \frac{\partial}{\partial A_{\zeta,e,i,l}} \int dv \mathbf{A} \cdot \mathbf{B} = & \\ -A_{\theta,e,j,p} \int ds T'_l T_p \oint d\theta d\zeta \cos \alpha_i \cos \alpha_j & \\ -A_{\theta,o,j,p} \int ds T'_l T_p \oint d\theta d\zeta \cos \alpha_i \sin \alpha_j & \\ +A_{\theta,e,j,p} \int ds T'_p T_l \oint d\theta d\zeta \cos \alpha_j \cos \alpha_i & \\ +A_{\theta,o,j,p} \int ds T'_p T_l \oint d\theta d\zeta \sin \alpha_j \cos \alpha_i & \end{aligned} \quad (17)$$

$$\begin{aligned} \frac{\partial}{\partial A_{\zeta,o,i,l}} \int dv \mathbf{B} \cdot \mathbf{B} = & \\ +2m_i(m_j A_{\zeta,o,j,p} + n_j A_{\theta,o,j,p}) & \int ds T_l T_p \oint d\theta d\zeta \bar{g}_{ss} \cos \alpha_i \cos \alpha_j \\ -2m_i(m_j A_{\zeta,e,j,p} + n_j A_{\theta,e,j,p}) & \int ds T_l T_p \oint d\theta d\zeta \bar{g}_{ss} \cos \alpha_i \sin \alpha_j \\ -2m_i A_{\zeta,e,j,p} & \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\theta} \cos \alpha_i \cos \alpha_j \\ -2m_i A_{\zeta,o,j,p} & \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\theta} \cos \alpha_i \sin \alpha_j \\ -2(m_j A_{\zeta,o,j,p} + n_j A_{\theta,o,j,p}) & \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\theta} \cos \alpha_j \sin \alpha_i \\ +2(m_j A_{\zeta,e,j,p} + n_j A_{\theta,e,j,p}) & \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\theta} \sin \alpha_j \sin \alpha_i \\ +2m_i A_{\theta,e,j,p} & \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\zeta} \cos \alpha_i \cos \alpha_j \\ +2m_i A_{\theta,o,j,p} & \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\zeta} \cos \alpha_i \sin \alpha_j \\ +2A_{\zeta,e,j,p} & \int ds T'_p T_l \oint d\theta d\zeta \bar{g}_{\theta\theta} \cos \alpha_j \sin \alpha_i \\ +2A_{\zeta,o,j,p} & \int ds T'_l T'_p \oint d\theta d\zeta \bar{g}_{\theta\theta} \sin \alpha_i \sin \alpha_j \\ -2A_{\theta,e,j,p} & \int ds T'_l T'_p \oint d\theta d\zeta \bar{g}_{\theta\zeta} \sin \alpha_i \cos \alpha_j \\ -2A_{\theta,o,j,p} & \int ds T'_l T'_p \oint d\theta d\zeta \bar{g}_{\theta\zeta} \sin \alpha_i \sin \alpha_j \end{aligned} \quad (18)$$

$$\begin{aligned} \frac{\partial}{\partial A_{\zeta,o,i,l}} \int dv \mathbf{A} \cdot \mathbf{B} = & \\ -A_{\theta,e,j,p} \int ds T'_l T_p \oint d\theta d\zeta \sin \alpha_i \cos \alpha_j & \\ -A_{\theta,o,j,p} \int ds T'_l T_p \oint d\theta d\zeta \sin \alpha_i \sin \alpha_j & \\ +A_{\theta,e,j,p} \int ds T'_p T_l \oint d\theta d\zeta \cos \alpha_j \sin \alpha_i & \\ +A_{\theta,o,j,p} \int ds T'_p T_l \oint d\theta d\zeta \sin \alpha_j \sin \alpha_i & \end{aligned} \quad (19)$$

1.9 second derivatives of energy and helicity integrands

1. The second derivatives wrt $A_{\theta,e,j,p}$ (stellarator symmetric) are

$$\begin{aligned} \frac{\partial}{\partial A_{\theta,e,j,p}} \frac{\partial}{\partial A_{\theta,e,i,l}} \int dv \mathbf{B} \cdot \mathbf{B} = & \\ + 2 n_i n_j \int ds T_l T_p \oint d\theta d\zeta \bar{g}_{ss} \sin \alpha_i \sin \alpha_j & \\ - 2 n_i \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\zeta} \sin \alpha_i \cos \alpha_j & \\ - 2 n_j \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\zeta} \sin \alpha_j \cos \alpha_i & \\ + 2 \int ds T'_l T'_p \oint d\theta d\zeta \bar{g}_{\zeta\zeta} \cos \alpha_i \cos \alpha_j & \end{aligned} \quad (20)$$

$$\begin{aligned} \frac{\partial}{\partial A_{\theta,o,j,p}} \frac{\partial}{\partial A_{\theta,e,i,l}} \int dv \mathbf{B} \cdot \mathbf{B} = & \\ - 2 n_i n_j \int ds T_p T_l \oint d\theta d\zeta \bar{g}_{ss} \cos \alpha_j \sin \alpha_i ! & \\ + 2 n_j \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\zeta} \cos \alpha_j \cos \alpha_i & \\ - 2 n_i \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\zeta} \sin \alpha_i \sin \alpha_j & \\ + 2 \int ds T'_l T'_p \oint d\theta d\zeta \bar{g}_{\zeta\zeta} \cos \alpha_i \sin \alpha_j ! & \end{aligned} \quad (21)$$

$$\begin{aligned} \frac{\partial}{\partial A_{\zeta,e,j,p}} \frac{\partial}{\partial A_{\theta,e,i,l}} \int dv \mathbf{B} \cdot \mathbf{B} = & \\ + 2 n_i m_j \int ds T_l T_p \oint d\theta d\zeta \bar{g}_{ss} \sin \alpha_i \sin \alpha_j & \\ + 2 n_i \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\theta} \sin \alpha_i \cos \alpha_j & \\ - 2 m_j \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\zeta} \sin \alpha_j \cos \alpha_i & \\ - 2 \int ds T'_p T'_l \oint d\theta d\zeta \bar{g}_{\theta\zeta} \cos \alpha_j \cos \alpha_i & \end{aligned} \quad (22)$$

$$\begin{aligned} \frac{\partial}{\partial A_{\zeta,o,j,p}} \frac{\partial}{\partial A_{\theta,e,i,l}} \int dv \mathbf{B} \cdot \mathbf{B} = & \\ - 2 n_i m_j \int ds T_p T_l \oint d\theta d\zeta \bar{g}_{ss} \cos \alpha_j \sin \alpha_i ! & \\ + 2 n_i \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\theta} \sin \alpha_i \sin \alpha_j & \\ + 2 m_j \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\zeta} \cos \alpha_j \cos \alpha_i & \\ - 2 \int ds T'_p T'_l \oint d\theta d\zeta \bar{g}_{\theta\zeta} \sin \alpha_j \cos \alpha_i & \end{aligned} \quad (23)$$

$$\frac{\partial}{\partial A_{\theta,e,j,p}} \frac{\partial}{\partial A_{\theta,e,i,l}} \int dv \mathbf{A} \cdot \mathbf{B} = 0 \quad (24)$$

$$\frac{\partial}{\partial A_{\theta,o,j,p}} \frac{\partial}{\partial A_{\theta,e,i,l}} \int dv \mathbf{A} \cdot \mathbf{B} = 0 \quad (25)$$

$$\begin{aligned} \frac{\partial}{\partial A_{\zeta,e,j,p}} \frac{\partial}{\partial A_{\theta,e,i,l}} \int dv \mathbf{A} \cdot \mathbf{B} = & \\ - \int ds T'_p T_l \oint d\theta d\zeta \cos \alpha_j \cos \alpha_i & \\ + \int ds T'_l T_p \oint d\theta d\zeta \cos \alpha_i \cos \alpha_j & \end{aligned} \quad (26)$$

$$\begin{aligned} \frac{\partial}{\partial A_{\zeta,o,j,p}} \frac{\partial}{\partial A_{\theta,e,i,l}} \int dv \mathbf{A} \cdot \mathbf{B} = & \\ - \int ds T'_p T_l \oint d\theta d\zeta \sin \alpha_j \cos \alpha_i & \\ + \int ds T'_l T_p \oint d\theta d\zeta \cos \alpha_i \sin \alpha_j ! & \end{aligned} \quad (27)$$

1.10 second derivatives of energy and helicity integrands

1. The second derivatives wrt $A_{\theta,e,j,p}$ (non-stellarator symmetric) are

$$\begin{aligned} \frac{\partial}{\partial A_{\theta,e,j,p}} \frac{\partial}{\partial A_{\theta,o,i,l}} \int dv \mathbf{B} \cdot \mathbf{B} = & \\ -2n_i n_j \int ds T_l T_p \oint d\theta d\zeta \bar{g}_{ss} \cos \alpha_i \sin \alpha_j ! & \\ +2n_i \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\zeta} \cos \alpha_i \cos \alpha_j & \\ -2n_j \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\zeta} \sin \alpha_j \sin \alpha_i & \\ +2 \int ds T'_p T'_l \oint d\theta d\zeta \bar{g}_{\zeta\zeta} \cos \alpha_j \sin \alpha_i ! & \end{aligned} \quad (28)$$

$$\begin{aligned} \frac{\partial}{\partial A_{\theta,o,j,p}} \frac{\partial}{\partial A_{\theta,o,i,l}} \int dv \mathbf{B} \cdot \mathbf{B} = & \\ +2n_i n_j \int ds T_l T_p \oint d\theta d\zeta \bar{g}_{ss} \cos \alpha_i \cos \alpha_j & \\ +2n_i \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\zeta} \cos \alpha_i \sin \alpha_j ! & \\ +2n_j \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\zeta} \cos \alpha_j \sin \alpha_i ! & \\ +2 \int ds T'_l T'_p \oint d\theta d\zeta \bar{g}_{\zeta\zeta} \sin \alpha_i \sin \alpha_j & \end{aligned} \quad (29)$$

$$\begin{aligned} \frac{\partial}{\partial A_{\zeta,e,j,p}} \frac{\partial}{\partial A_{\theta,o,i,l}} \int dv \mathbf{B} \cdot \mathbf{B} = & \\ -2n_i m_j \int ds T_l T_p \oint d\theta d\zeta \bar{g}_{ss} \cos \alpha_i \sin \alpha_j ! & \\ -2n_i \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\theta} \cos \alpha_i \cos \alpha_j & \\ -2m_j \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\zeta} \sin \alpha_j \sin \alpha_i & \\ -2 \int ds T'_p T'_l \oint d\theta d\zeta \bar{g}_{\theta\zeta} \cos \alpha_j \sin \alpha_i ! & \end{aligned} \quad (30)$$

$$\begin{aligned} \frac{\partial}{\partial A_{\zeta,o,j,p}} \frac{\partial}{\partial A_{\theta,o,i,l}} \int dv \mathbf{B} \cdot \mathbf{B} = & \\ +2n_i m_j \int ds T_l T_p \oint d\theta d\zeta \bar{g}_{ss} \cos \alpha_i \cos \alpha_j & \\ -2n_i \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\theta} \cos \alpha_i \sin \alpha_j ! & \\ +2m_j \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\zeta} \cos \alpha_j \sin \alpha_i ! & \\ -2 \int ds T'_p T'_l \oint d\theta d\zeta \bar{g}_{\theta\zeta} \sin \alpha_j \sin \alpha_i & \end{aligned} \quad (31)$$

$$\frac{\partial}{\partial A_{\theta,e,j,p}} \frac{\partial}{\partial A_{\theta,o,i,l}} \int dv \mathbf{A} \cdot \mathbf{B} = 0 \quad (32)$$

$$\frac{\partial}{\partial A_{\theta,o,j,p}} \frac{\partial}{\partial A_{\theta,o,i,l}} \int dv \mathbf{A} \cdot \mathbf{B} = 0 \quad (33)$$

$$\begin{aligned} \frac{\partial}{\partial A_{\zeta,e,j,p}} \frac{\partial}{\partial A_{\theta,o,i,l}} \int dv \mathbf{A} \cdot \mathbf{B} = & \\ - \int ds T'_p T_l \oint d\theta d\zeta \cos \alpha_j \sin \alpha_i ! & \\ + \int ds T'_l T_p \oint d\theta d\zeta \sin \alpha_i \cos \alpha_j & \end{aligned} \quad (34)$$

$$\begin{aligned} \frac{\partial}{\partial A_{\zeta,o,j,p}} \frac{\partial}{\partial A_{\theta,o,i,l}} \int dv \mathbf{A} \cdot \mathbf{B} = & \\ - \int ds T'_p T_l \oint d\theta d\zeta \sin \alpha_j \sin \alpha_i & \\ + \int ds T'_l T_p \oint d\theta d\zeta \sin \alpha_i \sin \alpha_j & \end{aligned} \quad (35)$$

1.11 second derivatives of energy and helicity integrands

1. The second derivatives wrt $\underline{A}_{\zeta,e,i,l}$ (stellarator symmetric) are

$$\begin{aligned} \frac{\partial}{\partial \underline{A}_{\theta,e,j,p}} \frac{\partial}{\partial \underline{A}_{\zeta,e,i,l}} \int dv \mathbf{B} \cdot \mathbf{B} = & \\ + 2 m_i n_j \int ds T_l T_p \oint d\theta d\zeta \bar{g}_{ss} \sin \alpha_i \sin \alpha_j & \\ + 2 n_j \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\theta} \sin \alpha_j \cos \alpha_i & \\ - 2 m_i \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\zeta} \sin \alpha_i \cos \alpha_j & \\ - 2 \int ds T'_l T'_p \oint d\theta d\zeta \bar{g}_{\theta\zeta} \cos \alpha_i \cos \alpha_j & \end{aligned} \quad (36)$$

$$\begin{aligned} \frac{\partial}{\partial \underline{A}_{\theta,o,j,p}} \frac{\partial}{\partial \underline{A}_{\zeta,e,i,l}} \int dv \mathbf{B} \cdot \mathbf{B} = & \\ - 2 m_i n_j \int ds T_p T_l \oint d\theta d\zeta \bar{g}_{ss} \cos \alpha_j \sin \alpha_i ! & \\ - 2 n_j \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\theta} \cos \alpha_j \cos \alpha_i & \\ - 2 m_i \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\zeta} \sin \alpha_i \sin \alpha_j & \\ - 2 \int ds T'_l T'_p \oint d\theta d\zeta \bar{g}_{\theta\zeta} \cos \alpha_i \sin \alpha_j ! & \end{aligned} \quad (37)$$

$$\begin{aligned} \frac{\partial}{\partial \underline{A}_{\zeta,e,j,p}} \frac{\partial}{\partial \underline{A}_{\zeta,e,i,l}} \int dv \mathbf{B} \cdot \mathbf{B} = & \\ + 2 m_i m_j \int ds T_l T_p \oint d\theta d\zeta \bar{g}_{ss} \sin \alpha_i \sin \alpha_j & \\ + 2 m_i \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\theta} \sin \alpha_i \cos \alpha_j & \\ + 2 m_j \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\theta} \sin \alpha_j \cos \alpha_i & \\ + 2 \int ds T'_l T'_p \oint d\theta d\zeta \bar{g}_{\theta\theta} \cos \alpha_i \cos \alpha_j & \end{aligned} \quad (38)$$

$$\begin{aligned} \frac{\partial}{\partial \underline{A}_{\zeta,o,j,p}} \frac{\partial}{\partial \underline{A}_{\zeta,e,i,l}} \int dv \mathbf{B} \cdot \mathbf{B} = & \\ - 2 m_i m_j \int ds T_p T_l \oint d\theta d\zeta \bar{g}_{ss} \cos \alpha_j \sin \alpha_i ! & \\ - 2 m_j \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\theta} \cos \alpha_j \cos \alpha_i & \\ + 2 m_i \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\theta} \sin \alpha_i \sin \alpha_j & \\ + 2 \int ds T'_l T'_p \oint d\theta d\zeta \bar{g}_{\theta\theta} \cos \alpha_i \sin \alpha_j ! & \end{aligned} \quad (39)$$

$$\begin{aligned} \frac{\partial}{\partial \underline{A}_{\theta,e,j,p}} \frac{\partial}{\partial \underline{A}_{\zeta,e,i,l}} \int dv \mathbf{A} \cdot \mathbf{B} = & \\ - \int ds T'_l T_p \oint d\theta d\zeta \cos \alpha_i \cos \alpha_j & \\ + \int ds T'_p T_l \oint d\theta d\zeta \cos \alpha_j \cos \alpha_i & \end{aligned} \quad (40)$$

$$\begin{aligned} \frac{\partial}{\partial \underline{A}_{\theta,o,j,p}} \frac{\partial}{\partial \underline{A}_{\zeta,e,i,l}} \int dv \mathbf{A} \cdot \mathbf{B} = & \\ - \int ds T'_l T_p \oint d\theta d\zeta \cos \alpha_i \sin \alpha_j ! & \\ + \int ds T'_p T_l \oint d\theta d\zeta \sin \alpha_j \cos \alpha_i & \end{aligned} \quad (41)$$

$$\frac{\partial}{\partial \underline{A}_{\zeta,e,j,p}} \frac{\partial}{\partial \underline{A}_{\zeta,e,i,l}} \int dv \mathbf{A} \cdot \mathbf{B} = 0 \quad (42)$$

$$\frac{\partial}{\partial \underline{A}_{\zeta,o,j,p}} \frac{\partial}{\partial \underline{A}_{\zeta,e,i,l}} \int dv \mathbf{A} \cdot \mathbf{B} = 0 \quad (43)$$

1.12 second derivatives of energy and helicity integrands

1. The second derivatives wrt $A_{\zeta,o,i,l}$ (non-stellarator symmetric) are

$$\begin{aligned} \frac{\partial}{\partial A_{\theta,e,j,p}} \frac{\partial}{\partial A_{\zeta,o,i,l}} \int dv \mathbf{B} \cdot \mathbf{B} = & \\ -2m_i n_j \int ds T_l T_p \oint d\theta d\zeta \bar{g}_{ss} \cos \alpha_i \sin \alpha_j & ! \\ +2n_j \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\theta} \sin \alpha_j \sin \alpha_i & \\ +2m_i \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\zeta} \cos \alpha_i \cos \alpha_j & \\ -2 \int ds T'_l T'_p \oint d\theta d\zeta \bar{g}_{\theta\zeta} \sin \alpha_i \cos \alpha_j & \end{aligned} \quad (44)$$

$$\begin{aligned} \frac{\partial}{\partial A_{\theta,o,j,p}} \frac{\partial}{\partial A_{\zeta,o,i,l}} \int dv \mathbf{B} \cdot \mathbf{B} = & \\ +2m_i n_j \int ds T_l T_p \oint d\theta d\zeta \bar{g}_{ss} \cos \alpha_i \cos \alpha_j & ! \\ -2n_j \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\theta} \cos \alpha_j \sin \alpha_i & \\ +2m_i \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\zeta} \cos \alpha_i \sin \alpha_j & ! \\ -2 \int ds T'_l T'_p \oint d\theta d\zeta \bar{g}_{\theta\zeta} \sin \alpha_i \sin \alpha_j & \end{aligned} \quad (45)$$

$$\begin{aligned} \frac{\partial}{\partial A_{\zeta,e,j,p}} \frac{\partial}{\partial A_{\zeta,o,i,l}} \int dv \mathbf{B} \cdot \mathbf{B} = & \\ -2m_i m_j \int ds T_l T_p \oint d\theta d\zeta \bar{g}_{ss} \cos \alpha_i \sin \alpha_j & ! \\ -2m_i \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\theta} \cos \alpha_i \cos \alpha_j & \\ +2m_j \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\theta} \sin \alpha_j \sin \alpha_i & \\ +2 \int ds T'_p T'_l \oint d\theta d\zeta \bar{g}_{\theta\theta} \cos \alpha_j \sin \alpha_i & ! \end{aligned} \quad (46)$$

$$\begin{aligned} \frac{\partial}{\partial A_{\zeta,o,j,p}} \frac{\partial}{\partial A_{\zeta,o,i,l}} \int dv \mathbf{B} \cdot \mathbf{B} = & \\ +2m_i m_j \int ds T_l T_p \oint d\theta d\zeta \bar{g}_{ss} \cos \alpha_i \cos \alpha_j & ! \\ -2m_i \int ds T_l T'_p \oint d\theta d\zeta \bar{g}_{s\theta} \cos \alpha_i \sin \alpha_j & ! \\ -2m_j \int ds T_p T'_l \oint d\theta d\zeta \bar{g}_{s\theta} \cos \alpha_j \sin \alpha_i & ! \\ +2 \int ds T'_l T'_p \oint d\theta d\zeta \bar{g}_{\theta\theta} \sin \alpha_i \sin \alpha_j & ! \end{aligned} \quad (47)$$

$$\begin{aligned} \frac{\partial}{\partial A_{\theta,e,j,p}} \frac{\partial}{\partial A_{\zeta,o,i,l}} \int dv \mathbf{A} \cdot \mathbf{B} = & \\ - \int ds T'_l T_p \oint d\theta d\zeta \sin \alpha_i \cos \alpha_j & \\ + \int ds T'_p T_l \oint d\theta d\zeta \cos \alpha_j \sin \alpha_i & ! \end{aligned} \quad (48)$$

$$\begin{aligned} \frac{\partial}{\partial A_{\theta,o,j,p}} \frac{\partial}{\partial A_{\zeta,o,i,l}} \int dv \mathbf{A} \cdot \mathbf{B} = & \\ - \int ds T'_l T_p \oint d\theta d\zeta \sin \alpha_i \sin \alpha_j & \\ + \int ds T'_p T_l \oint d\theta d\zeta \sin \alpha_j \sin \alpha_i & ! \end{aligned} \quad (49)$$

$$\frac{\partial}{\partial A_{\zeta,e,j,p}} \frac{\partial}{\partial A_{\zeta,o,i,l}} \int dv \mathbf{A} \cdot \mathbf{B} = 0 \quad (50)$$

$$\frac{\partial}{\partial A_{\zeta,o,j,p}} \frac{\partial}{\partial A_{\zeta,o,i,l}} \int dv \mathbf{A} \cdot \mathbf{B} = 0 \quad (51)$$

1.13 matrix elements

1. The energy, $W \equiv \int dv \mathbf{B} \cdot \mathbf{B}$, and helicity, $K \equiv \int dv \mathbf{A} \cdot \mathbf{B}$, functionals may be written

$$W = \frac{1}{2} a_i A_{i,j} a_j + a_i B_{i,j} \psi_j + \frac{1}{2} \psi_i C_{i,j} \psi_j \quad (52)$$

$$K = \frac{1}{2} a_i D_{i,j} a_j + a_i E_{i,j} \psi_j + \frac{1}{2} \psi_i F_{i,j} \psi_j \quad (53)$$

where $\mathbf{a} \equiv \{A_{\theta,e,i,l}, A_{\zeta,e,i,l}, A_{\theta,o,i,l}, A_{\zeta,o,i,l}, f_{e,i}, f_{o,i}\}$ contains the independent degrees of freedom and $\psi \equiv \{\Delta\psi_t, \Delta\psi_p\}$.

2. The matrix elements are computed via

$$\text{MA(i,j)} \equiv A_{i,j} = \frac{\partial^2 W}{\partial a_i \partial a_j} \quad (54)$$

$$\text{MB(i,j)} \equiv B_{i,j} = \frac{\partial^2 W}{\partial a_i \partial \psi_j} \quad (55)$$

$$\text{MC(i,j)} \equiv C_{i,j} = \frac{\partial^2 W}{\partial \psi_i \partial \psi_j} \quad (56)$$

3. The energy functionals can also be represented as

$$\begin{aligned} W &= \frac{1}{2} \mathbf{a}^T \cdot A[\mathbf{x}] \cdot \mathbf{a} + B[\mathbf{x}, \psi] \cdot \mathbf{a} + C[\mathbf{x}, \psi], \\ K &= \frac{1}{2} \mathbf{a}^T \cdot D \cdot \mathbf{a} + E[\psi] \cdot \mathbf{a} + F[\psi]. \end{aligned} \quad (57)$$